

THERMAL PARAMETER ESTIMATION USING RECURSIVE IDENTIFICATION

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Abstract: A novel method that converts a semiconductor Transient Thermal Impedance Curve (TTIC) into an equivalent thermal R - C network model is presented. Thermal Resistance (R) and thermal Capacitance (C) parameters of the model are identified using manufacturer's data and off line Recursive Least Square (RLS) techniques. Relevant estimation theory concepts and the formulation of an appropriate model for the identification process are given. Model synthesis is illustrated using an isolated base power transistor module. The application of time decoupled theory for high order thermal models is outlined. Simulation of junction temperature responses using model and manufacturer TTIC's are compared. Identified parameter validity is further confirmed by parameter calculation obtained from module physical dimensions.

In this paper, a new approach to the problem of determining device thermal characteristics is presented from the System Identification point of view [9]. Sections 2 and 3 outline the basic identification procedure used and the governing laws for device thermal model building. Sections 4 and 5 review and apply estimation theory principles to determine the R - C parameter values given the manufacturer's TTIC curve. Section 6 presents identification results for an isolated base transistor model example and discusses parameter accuracy of the identified R - C values.

I. Introduction

Power converter manufacturers typically utilize the relatively low but significant thermal heat capacity of power semiconductors to obtain short duration overload ratings well in excess of continuous ratings. The methods for determining the peak allowable junction temperature ($T_{j(max)}$) under transient and intermittent loading are well established and have remained essentially unchanged since 1959 [1]. The standard approach uses the TTIC supplied by device manufacturers (Fig. 1a). Junction temperature response to device power pulses are estimated from this curve and the principle of superposition for conditions such as single pulse overload, repetitive pulse overload, overloads following continuous duty and irregularly shaped power vs. time profiles. Indeed, power transistor Safe Operating Area (SOA) limits are usually based upon this approach [2-5]. However, transient junction temperature estimation using the TTIC approach has several shortcomings.

(i) *a-posteriori calculations* - Circuit simulation programs containing sophisticated device models exist that can calculate instantaneous power vs. time profiles [6-7]. However, instantaneous junction temperature vs. time profiles cannot be solved with the TTIC concept until the entire overload simulation process is complete, stored in memory and broken down into equivalent pulse amplitude and durations. This inefficiency suggests the use of a device thermal model (Fig. 1b & 1c.) for maximizing simulation capability by solving both the electrical and thermal network models simultaneously.

(ii) *graphical analysis*- The standard TTIC approach requires cumbersome graphical analysis to transform the irregularly shaped power profiles, such as the switching and conduction loss profiles in SOA calculations, into equivalent energy [watt-sec] "square" power pulses upon which this curve is based.

(iii) *repeat calculations* - Each application of a new overload sequence requires T_j to be recalculated.

(iv) *desired accuracy* - The accuracy of estimated junction temperatures decreases for increasingly complex overload waveforms such as Pulse Width Modulation (PWM) acceleration of a motor. When using the standard approach, gross simplifying assumptions are necessary to keep graphical analysis and hand computations tractable.

The shortcomings of the standard approach suggest the need to develop an accurate thermal model to make better estimates of T_j . Use of a device thermal model for indirect measurement of the junction to case temperature rise (ΔT_{jc}) may result in improved converter fault diagnostics. Indirectly calculating ΔT_{jc} in real time may be done with the discrete transfer function model of Fig. 1b and a Digital Signal Processor (DSP) / microprocessor or an analog operational amplifier model of Fig. 1c. In addition, potential advantages in substantially increased converter overload ratings exist when using the ΔT_{jc} "observer" based model in real time adaptive control [8].

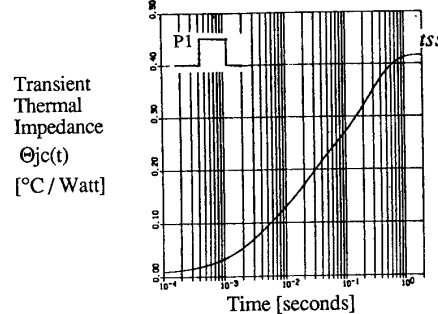


Fig. 1a Semiconductor Transient Thermal Impedance Curve (junction-case)

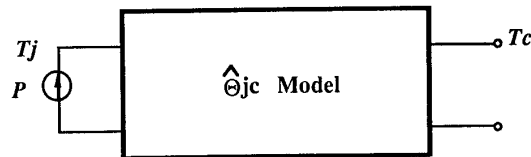


Fig. 1b Estimated Black Box Model of Fig. 1a

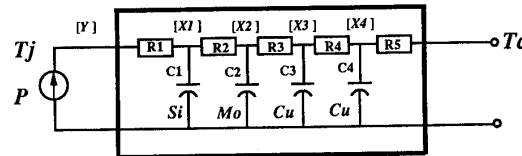


Fig. 1c Thermal Model Structure of Fig. 1a & Fig. 1b

II. Identification Procedure

The success or failure of any application of system identification rests on finding an algorithm that can intelligently utilize apriori information. The thermal estimation problem has three main features which can be exploited. First, the structure of Fig. 1c is known from physical grounds to closely model the thermal behavior of the system, even though the exact values of the R 's and C 's are unknown. This suggests that one of the parametric identification methods should be applicable. The second feature is the TTIC curve,

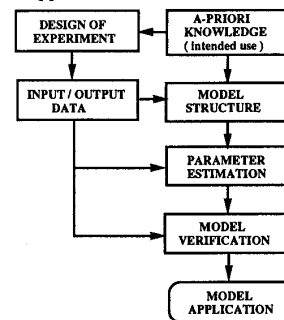


Fig. 2 Identification Procedure

which can be used to construct simulated input/output pairs to sufficiently excite the identification procedure. The third feature is that the thermal time constants have a wide time scale separation. In many identification setups, this would be a serious problem because it is difficult to excite all modes of the system without an inordinately large number of time steps. Since this time scale separation is known to exist a priori, however, we are able to exploit it by identifying the slow and fast modes separately via a time decoupling approach. Figure 2 defines the 4 basic steps used [10].

(1) **Model Formulation** : The type and order of the thermal model structure are defined from a-priori knowledge about the semiconductor. Some numerical constants of the model can also be obtained a-priori by applying Newtons Law of Cooling. Discrete state space equations are derived based upon the physical model.

(2) **Design of Experiment** : The input signal, sampling interval and experiment length are chosen so that appropriate modes of the thermal model are sufficiently excited for identification.

(3) **Parameter Estimation** : This step determines the numerical values of the model structure. The choice of algorithm is the off-line direct method of Recursive Least Squares (RLS). The method shown in Fig. 3 is based upon the fact that the collected output responses are linearly dependent on the unknown parameters where;

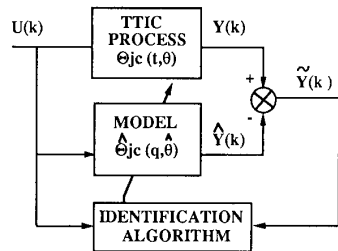


Fig. 3 Parameter Estimation Technique

k = Discrete sampling time

$u(k)$ = Input power pulse sequence into the device.

$y(k)$ = Calculated junction-case temperature response, $T_{jc}(k)$

$\hat{y}(k)$ = Estimated junction-case temperature response, $\hat{T}_{jc}(k)$

$\tilde{y}(k)$ = Output error equation between true $T_{jc}(k)$ and estimated $\hat{T}_{jc}(k)$ responses.

The basic scheme is to use a well planned input power sequence which has a sampling interval shorter than the fastest time constant and an experiment length that is longer than the slowest estimated time constant of the thermal model. True junction temperature response, $T_{jc}(k)$, to $u(k)$ is calculated using the manufacturer's TTIC. The same input test signal is also applied, restarting at $t = 0$, to the $\hat{\theta}_{jc}$ model and $\hat{y}(k)$ is calculated for each time instant k . The RLS algorithm attempts to drive the $\tilde{y}(k)$ error to zero at each k by adapting the unknown parameters $\hat{\theta}_{jc}$. The estimated numerator and denominator coefficients of $\hat{\theta}_{jc}$ converge to steady state values for a properly designed input test signal. The actual Resistance and Capacitance values may then be determined from these coefficients using the initial model equations formulated.

(4) **Model Verification** : This step relates the identification (ID) results to well known physical results. Comparison of $T_{jc}(k)$ vs. $\hat{T}_{jc}(k)$ and Manufacturer's TTIC vs. an estimated TTIC curve are made. Additionally, identified R and C model parameters are compared to calculated $R-C$ values obtained from measurements of an actual semiconductor. These 4 steps are now examined in more detail.

III. Model Formulation

Transient Thermal Impedance Curve:

The semiconductor thermal model structure is implicitly contained in the TTIC as a complex sum of $R-C$ exponentials. It is therefore

desirable to review the definition, derivation, assumptions and application of this curve. The concept of thermal resistance is based upon an analogy between electrical and thermal systems with temperature[°C], heat flow due to power dissipation [Watts] and thermal resistance [°C/W] being analogous to voltage, current and electrical resistance [11]. The TTIC in Fig. 1a is obtained by applying a single "square" power pulse "P1" to the device until the junction temperature reaches steady state at time t_{ss} . Junction temperature rise $\{\Delta T_{jc}\}$ is determined by fixing the case constant at ambient temperature $\{T_c=T_a\}$ and measuring device temperature with infrared methods or electrical Temperature Sensitive Parameters $\{TSP\}$ such as forward voltage drop $\{V_f\}$ or base emitter voltage $\{V_{be}\}$ [12]. The actual ΔT_{jc} rise is found by correlating the measured change in V_{be} ($\cong 2\text{mv} / ^\circ\text{C} / \text{junction}$) vs. time to a previous calibrated V_{be} vs Temperature test for constant T_a and base current. The transient thermal impedance is defined at any time t as

$$\Theta_{jc}(t) = \frac{T_j(t) - T_c(t)}{P_1} = \frac{\Delta T_{jc}(t)}{P_1} \quad (1)$$

The thermal system is assumed to be linear, and hence superposition can be applied to the TTIC. The TTIC is a "step response" curve with zero initial conditions, relating device step input power to ΔT_{jc} at the

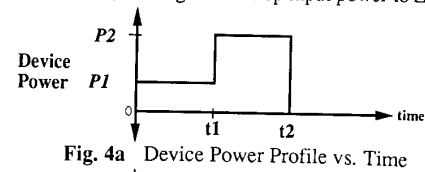


Fig. 4a Device Power Profile vs. Time

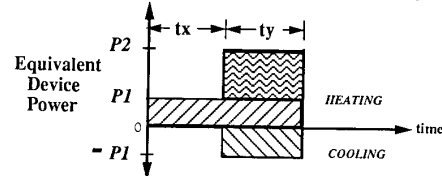


Fig. 4b Equivalent Heating and Cooling Pulses

output. The power profile in figure 4a can be separated into equivalent heating and cooling pulse durations of t_x and t_y as shown in Fig. 4b. The junction temperature rise can be determined by adding individual T_{jc} pulse responses [1].

$$\Delta T_{jc}(t_1) = P_1 * \Theta_{jc}(t_1) \quad (2)$$

$$\Delta T_{jc}(t_2) = P_1 * \Theta_{jc}(t_2) - P_1 * \Theta_{jc}(t_2-t_1) + P_2 * \Theta_{jc}(t_2-t_1) \quad (3)$$

The published TTIC is usually higher than the tested value to account for manufacturing variations and the increase in thermal resistance over time.

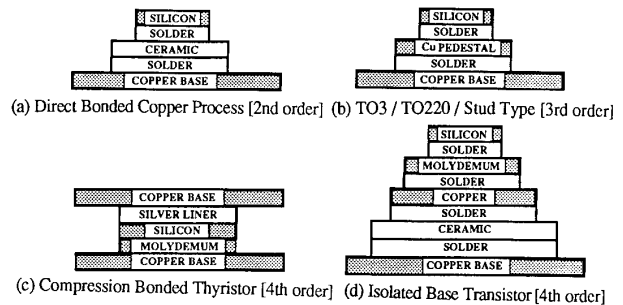


Fig. 5 Model Order Dependent on Semiconductor Package

A-priori Knowledge

Knowledge of the physical properties of the semiconductor can be used to fix the model structure and order, and to determine some numerical values of the "Black Box" shown in Fig. 1c. The use of all available a-priori knowledge prior to application of the identification algorithm is important since misleading results due to an assumed

wrong structure are difficult to detect from data alone. Also, a-priori knowledge can enhance model validity and model accuracy. An appropriate model can be formulated using (i) physical knowledge, (ii) I/O measurements, or both.

(i) Physical Knowledge: The model order is dependent on the type of semiconductor package used as shown in Fig. 5. The exact order can be determined by visual inspection of the package cross sectional view and replacing significant heat capacity materials (Cu, Si, Mo) with thermal capacitances. The number of capacitors determines the model order. One dimensional heat flow from junction to case results in the typical R-C network structure shown in Fig. 1c. Numerical values for $R5$ and $C4$ of the copper base can be obtained without package dis-assembly by applying the governing thermal laws defined in Appendix 1 to a 4th order 50 Amp isolated base transistor.

$$R_i = L_i / (K * Ae) \quad [^{\circ}C / W] \quad (4)$$

$$C_i = \rho * Cp * V \quad [W-S / ^{\circ}C] \quad (5)$$

Thermal resistance and capacitance calculations can be extended to $R1-R4$ and $C1-C3$ by dis-assembling the package and physically measuring each material thickness and cross section area perpendicular to heat flow. This "Calculated Parameter" approach is documented in Appendix A1 with the results shown in Table 1. An analog simulation of this R-C structure is plotted in Fig. 6 vs. the actual TTIC. This procedure alone may produce sufficient $\Theta_{jc}(t)$ accuracy for the intended use of the model.

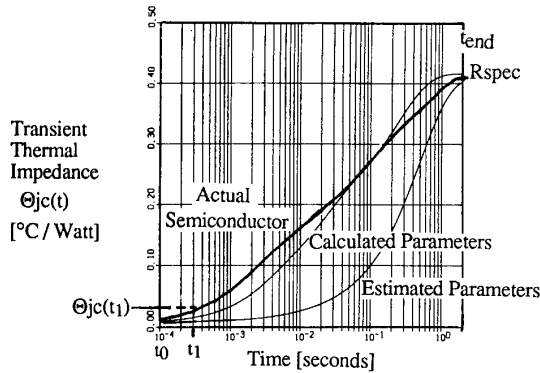


Fig. 6 Transient Thermal Impedance Curves

(ii) I/O Measurement: The asymptotic behaviour at the origin and steady state time t_{ss} of the manufacturer's TTIC in Fig. 6 can provide numerical values for parameters $R1, C1$ and for the sum of $R1$ thru $R5$ in Fig. 1c. The horizontal asymptote at the origin reflects the silicon thermal resistance $R1$. The initial slope near the origin of Fig. 6 can be \approx silicon thermal capacity by the analogy to $i = C / (dV/dT)$:

$$Power \approx C1 \Delta T_{jc} [^{\circ}C] / \Delta Time [sec] \quad [Watts] \quad (6)$$

$$C1 \approx P * (t1 - t0) / P * (\Theta(t1) - \Theta(t0)) \quad [W-S / ^{\circ}C] \quad (7)$$

The sum of $R1-R5$ is the "DC gain" of the mathematical transfer function model and is the value $\Theta_{jc}(t_{ss})$ read from the TTIC.

$\hat{\Theta}_i$ Parameter Initialization Procedure

As an alternative to the physical R-C calculation procedure above, the following procedure may be used to find initial estimates for the RLS identification algorithm using only TTIC information. Although this method involves crude approximations, it is helpful in estimating the model time constants for selecting a suitable RLS sampling time and providing parameters sufficiently close to the actual so that the RLS routine converges rapidly.

The following known data in Eq (8a-8d) can be obtained from the manufacturer's data specification sheet, the TTIC of Fig. 6 or from external physical dimensions.

$$R_{spec} = (R1 + R2 + R3 + R4) + R5 \quad (8a)$$

$$R5 \text{ is calculable from baseplate case thickness dimensions} \quad (8b)$$

$$R1 \text{ is from the TTIC as in the previous section} \quad (8c)$$

$$\tau_4 \text{ is from the TTIC since } t_{end} \text{ is known and is } 5\tau_4 \quad (8d)$$

Crude approximations for a $n = 4$ th order system are:

$$t_{end} \approx [5\tau_4] \approx 5[(R1 + R2 + R3 + R4) / C4] \quad (8e)$$

$$\tau_4 \approx [5\tau_3] \approx 5[(R1 + R2 + R3) / C3] \quad (8f)$$

$$\tau_3 \approx [5\tau_2] \approx 5[(R1 + R2) / C2] \quad (8g)$$

$$\tau_2 \approx [5\tau_1] \approx 5[(R1 * C1)] \quad (8h)$$

$$\tau_1 \approx [R1 * C1] \approx [t_{end} / 5^n] \quad (8i)$$

Thermal capacitance is assumed to be increasing by a constant factor MF_i from one stage to the next.

$$C2 \approx MF1 * C1 \quad (8j) \quad C3 \approx MF2 * C2 \quad (8k)$$

$$C4 \approx MF3 * C3 \quad (8l) \quad C4 / C1 \approx MF1 * MF2 * MF3 \quad (8m)$$

$$MF_i \approx (C4 / C1)^{1/3} \quad (8n)$$

Parameter $C4$ is found using Eq(a, b, d & e). Capacitor $C1$ is found using Eq(c & i). Capacitors $C2$ & $C3$ are extracted from Eq(j, k, & n). Resistor $R4$ is found by solving for the $\{R1+R2+R3\}$ sum of Eq(f) and substituting into Eq(e). Resistor $R3$ is found by solving for the sum $\{R1+R2\}$ of Eq(g) and substituting into Eq(f). Resistor $R2$ is calculated directly from Eq(g). Table 1 summarizes the Estimated R-C parameter results for the known data shown below.

$$t_{end} = 1.8 \text{ sec.} \quad R_{spec} = 0.41 [^{\circ}C/W] \quad R1 = 0.008 [^{\circ}C/W] \quad R5 = 0.0106 [^{\circ}C/W]$$

Table 1 Calculated & Estimated R-C Parameters

	$R1$	$R2$	$R3$	$R4$	$R5$	$C1$	$C2$	$C3$	$C4$
Calc.	0.006	0.111	0.122	0.166	0.011	0.033	0.148	1.180	9.500
Est.	0.008	0.022	0.078	0.291	0.011	0.360	0.488	0.665	9.000

Assumptions

The thermal model for Fig. 1 and Fig. 5 is assumed to be a linear, n th order lumped parameter, time invariant, deterministic, Single Input Single Output (SISO) system. Nonlinear radiation effects which are proportional to the 4th power of temperature are not significant since one dimensional heat flow is mostly by conduction. However, silicon conductivity is nonlinear with temperature varying 2:1 over the 25-150 $^{\circ}C$ operating range and may effect the estimates of the $R1$ & $R2$ thermal resistances. Also, some insulators such as BeO will vary by 20% over the same range. The present procedure ignores these nonlinearities, though their incorporation into the design procedure is an important area for further investigation. Lastly, measurement noises are assumed negligible.

Equations

A thermal model can be formulated using either an internal state space or I/O transfer function model approach. The first method is the most desirable since it is directly related to the physical structure of Fig. 1c. However, the need to measure the internal states (x_1, x_2, x_3, x_4) to find model coefficients precludes its use. The transfer function approach must be used since only I/O data from the TTIC is available. The disadvantage of this approach is that the Θ_{jc} model coefficients obtained from I/O data have no direct physical meaning. The R-C parameters of the structure are hidden in the numerator and denominator coefficients and must be further extracted. The transfer function model is developed in the continuous time domain and must be further discretized for use in the identification algorithm.

Continuous State Space Model

The system equations for the 4th order structure in Fig. 1c can be obtained using the capacitor voltages (temperatures) as the state variables (x_1, x_2, x_3, x_4). The output equation variable (Y) represents the silicon absolute junction temperature T_j . The input variable (u) represents the device power in watts. The transfer function Θ_{jc} reflects the temperature rise for a given power input. It is derived by applying the Laplace transform operator $\{s\}$ to

Eq(9a), solving for X and substituting into Eq(10a).

$$\dot{X} = AX + Bu \quad [^{\circ}C/\Delta t] \quad (9a)$$

$$Y = CX + Du \quad [^{\circ}C] \quad (10a)$$

$$Y/u = C(sI - A)^{-1}B + D \quad [^{\circ}C/W] \quad (11a)$$

where $B = (\frac{1}{C1}, 0, 0, 0)^T$, $C = (1, 0, 0, 0)^T$, $D = R1$, and

$$A = \begin{pmatrix} \frac{-1}{R2C1} & \frac{1}{R2C1} & 0 & 0 \\ \frac{-1}{R2C2} & \left(\frac{-1}{R2C2} - \frac{1}{R3C2} \right) & \frac{1}{R3C2} & 0 \\ 0 & \frac{1}{R3C3} & \left(\frac{-1}{R3C3} - \frac{1}{R4C3} \right) & \frac{1}{R4C3} \\ 0 & 0 & \frac{-1}{R4C4} & \left(\frac{-1}{R4C4} - \frac{1}{R5C4} \right) \end{pmatrix}$$

The symbolic transfer function corresponding to Eq(11) contains 4 numerator and 4 denominator coefficients in the s^0 to s^4 powers. Determination of the R - C parameter values requires simultaneously solving the 8 coefficient equations. Four of the 8 coefficients each contain 21 nonlinear sum and product terms of the form $1/(R_x C_x)$, which makes solution a formidable task. Using the apriori knowledge that the values of the capacitors are widely separated in this 4th order thermal model, the above equation can be split into two cascaded 2nd order systems which match the dc gain and overall dynamics of the original system. Thus, only simpler 2nd order equations need be developed using $x_1, x_2, R1, R2, R3, C1$ and $C2$.

$$\begin{pmatrix} \dot{x}_1 \\ \dot{x}_2 \end{pmatrix} = \begin{pmatrix} \frac{-1}{R2C1} & \frac{1}{R2C1} \\ \frac{-1}{R2C2} & \left(\frac{-1}{R2C2} - \frac{1}{R3C2} \right) \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} + \begin{pmatrix} \frac{1}{C1} \\ 0 \end{pmatrix} P \quad (12)$$

$$Y = [\quad 1 \quad 0 \quad] X^T + [R1] P \quad (13)$$

This is the time scale decoupling, and is possible whenever the unknown system has widely separated modes. In essence, the identification of the slow modes is conducted separately from the identification of the fast modes.

Discrete Transfer Function Model

In order to utilize an appropriate identification algorithm, the continuous system (12) and (13) must be discretized for computer implementation. There are several possible methods such as Eulers methods, Tustins approximation, step invariance, etc. An infinite series approximation was chosen because it leads to relatively simple equations relating the RC parameters to the filter coefficients. The discrete system equations are defined using the time shift operator, $\{q\}$.

$$X(k+1) = q X(k) = \Phi X(k) + \Gamma u(k) \quad [^{\circ}C/\Delta t] \quad (14)$$

$$Y(k) = C X(k) + D u(k) \quad [^{\circ}C] \quad (15)$$

The pulse transfer function is derived by solving Eq(14) for $x(k)$ and substituting into Eq(15) [13].

$$\hat{\Theta}_{jc}(k) = Y(k)/u(k) = C(qI - \Phi)^{-1}\Gamma + D \quad [^{\circ}C/W] \quad (16)$$

where h = sample interval

Φ = state transition matrix

$\Phi = e^{Ah} = I + Ah$ { first order approximation }

$$\Gamma = \int_0^h e^{A\zeta} d\zeta B \quad (17)$$

Performing the matrix operations and applying the backward shift operator (q^{-1}) yields the standard digital filter format with the coefficients defined in Appendix A2.

$$\hat{\Theta}_{jc}(q, \hat{\theta}) = \hat{H}(q, \hat{\theta}) = \frac{(b_0 + b_1 q^{-1} + b_2 q^{-2})}{(1 + a_1 q^{-1} + a_2 q^{-2})} \quad (18)$$

where $\hat{\theta} = [a_1, a_2, b_0, b_1, b_2]^T$
 n = order of the numerator = 2
 m = order of the denominator = 2

The linear difference equation for the system is;

$$y(k) = a_1 y(k-1) + a_2 y(k-2) + b_0 u(k) + b_1 u(k-1) + b_2 u(k-2) \quad (19)$$

Discretization Error

Substitution of $q = 1$ into Eq(18) provides an estimate for the discretization error due to a first order approximation for the state transition matrix.

$$\hat{\Theta}_{jc}(1, \hat{\theta}) = \hat{H}(1, \hat{\theta}) = \|R_1 + R_2 + R_3\| - \{0.5h / C1\} \quad (20)$$

The first term is the correct DC gain of the 2nd order system while the second term is due to the $O(h^3)$ error defined in Appendix A2. This term is negligible for small sample times used in identifying the fast time constants but can lead to large dc gain errors ($\approx 30\%$) in $y(k)$ for large sample time $\{h\}$. This is further detailed in Section VI.

Parameter Error Estimation

A worst case steady state error bound estimate for parameter R_1 due to Eq(20) can be derived using the standard rule of thumb that the system is sampled 10x faster than the fastest time constant to be identified.

$$\{h * 0.5 / C1\} = \left(\frac{\tau_{fastest}}{10} \right) \left(\frac{0.5}{C1} \right) = \left(\frac{R_1 C_1}{20 C_1} \right) = \pm 5\% \|R_1\| \quad (21)$$

R-C Parameter Extraction from Filter Coefficients

The five R - C parameters may be found by simultaneously solving the 5 coefficient equations in Appendix A2. A convergent solution is obtained if the $O(h^3)$ error term in the b_2 equation is eliminated before solving. The parameter solution equations [22-26] must be solved in the sequential order as shown. These equations were solved by hand using the variable substitution method and verified using the symbolic equation solver *Mathematica* [14].

$$R_1 = b_0 \quad (22)$$

$$C_1 = h a_1 / \{b_0 [(a_1 - a_2) + (b_2 - b_1) / b_0]\} \quad (23)$$

$$R_2 = (-0.5) h^2 / \{b_0 C_1^2 [(b_1 / b_0) - a_1] - C_1 h\} \quad (24)$$

$$C_2 = h / \{R_2 [\{a_1 + 2.0 - (h/R_2 C_1)\} - (a_1 + a_2 + 1.0)(R_2 C_1 / h)]\} \quad (25)$$

$$R_3 = h^2 / \{C_1 C_2 R_2 (a_1 + a_2 + 1.0)\} \quad (26)$$

IV. Design Of The Experiment

Standard Identification Procedure

In the semiconductor thermal model structure of Fig. 1c, the time constants (τ_i) cover a 1500:1 range from μ sec to seconds. Identifying these major time constants by standard identification methodology requires multiple "trial & error" ID applications, since coefficient accuracy tends to degenerate for systems with τ 's having more than 2 decades of time separation. Such experiments to identify the τ 's require engineering tradeoffs regarding sampling time $\{Ts\}$, experiment record length, and input signal amplitude. Proper identification of the fast time constant requires a high sampling rate, identification of the slow time constant requires a long record of input/output data. Together, these imply a cumbersome and poorly conditioned identification setup. One approach is to collect a number of experiments and to average them to obtain an averaged transfer function model that drives the $y(k) - \hat{y}(k)$ error to zero for a specific time region of interest. Further, ambiguous sets of R - C parameters may result if the R - C Extraction procedure of Section III is applied to such "averaged" models.

Proposed Time Decoupled Theory (TDT)

The disadvantages of widely separated time constants can be circumvented since we know apriori that such a separation exists.

The basic strategy of time decoupled identification is to run two "separate" identification procedures, one for the slow modes and one for the fast modes. Besides the advantages of tailoring the sampling rates and record lengths to the expected order of magnitude of the time constant, this decoupling allows a simpler R-C parameter extraction procedure. Figure 7 conveys the TDT procedure as applied to the 4th order semiconductor module example. The 4th order model is decoupled into two independent 2nd order systems. The split model is reasonable given the analogous electrical model where a high frequency device power sequence of short experiment length will charge $C1$ and $C2$ while leaving $C3$ and $C4$ virtually unchanged. Similarly, the fast modes will be virtually invisible to a step inputs with a slow sampling rate. The basic TDT concept involves using multiple identification runs to estimate the 4 major time constants in succession from the fastest to the slowest.

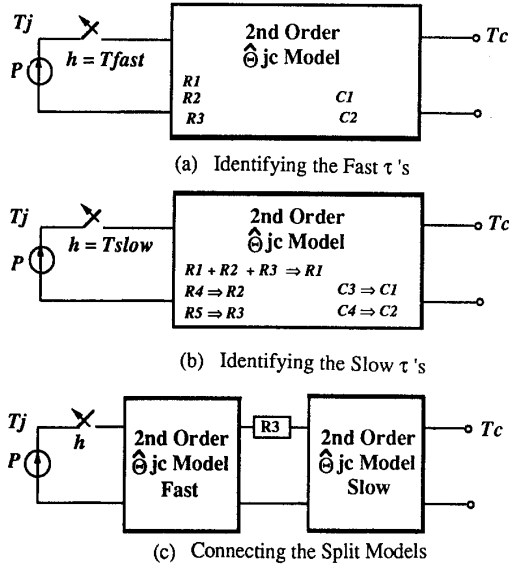


Fig. 7 Time Decoupled Theory Procedure (TDT)

The fastest time constant $[\tau_1]$ of Fig. 7a is identified by suitable selection of (i) a sampling rate that is fast enough for the estimated $[\tau_1]$ to be identified, (ii) a persistently exciting device power sequence, (iii) an experiment record length long enough to allow for parameter convergence and (iv) using all available a-priori knowledge for $R1, R2, R3, C1$ and $C2$ initial estimates. True junction temperature $y(k)$ is calculated using the TTIC. The identified a_i and b_i coefficients will typically converge in less than $k = 15$ samples. The $\hat{\theta}_i$ parameters are then passed to the R-C extraction procedure. The newly updated values for $R1$ and $C1$ will be very close to the actual values while $R2, R3$ and $C2$ values will be relatively inaccurate.

The ID procedure is next repeated using the R-C parameters from the previous run as initial a-priori estimates. The sampling rate is now chosen to be greater than Ts_1 but still fast enough to identify the second guess-timed time constant τ_2 . Application of the ID algorithm and parameter extraction procedures will modify $R2, R3$ and $C2$ to the correct actual values while the faster $R1-C1$ parameters will remain unchanged.

To identify the slower time constants $[\tau_3 \& \tau_4]$, the split model of Fig. 7b is used with the same second order equations as was used in Fig. 7a. A key parameter change is the substitution of $\{R1+R2+R3\} \Rightarrow R1$ to maintain the TTIC overall DC gain and slower system dynamics. The sampling rate is selectively increased and the ID & R-C extraction procedures are similarly repeated to find $R4, R5, C3$ and $C4$ in about 2 or 3 ID runs. Some systems may require repeating this fast / slow identification procedure to more accurately identify the interconnecting $R3$ element in Fig. 7c. Considerations for choosing a suitable sampling rate, experiment length and input signal amplitude are now discussed in more detail.

Input Signal

The amplitude of the device input power signal should be as high as allowable to improve accuracy. The form of the input signal should (1) consist of square pulses so that the TTIC curve for $y(k)$ calculation may be directly used. (2) have a random amplitude vs. time profile to allow ID convergence to a unique set of parameter values. (3) never have a non-realistic negative power pulse. (4) ideally result in rated T_j for rated power $\{Prated\}$ with steady state thermal resistance at $T_c = 25$ C. These constraints are met by taking a pseudo random sequence consisting of the first positive 50 digits of pi (0-9). For long sequences, the "average" random digit value approaches 5. The $u(k)$ amplitude equation used is:

$$u(k) = 2.0 * Prated * \{Random\ digit\} / 10 \quad (27)$$

To help identify longer term dynamics, six similar amplitude pulses $\{eg., u(1)$ to $u(6)\}$ were grouped together before the next allowable amplitude change.

Sampling Time

The Nyquist theorem determines the minimum sampling rate to use for each ID run. A commonly used practical rule of thumb is to sample 10x faster than the fastest time constant to be identified.

$$h = \tau_{smallest} / 10 \quad (28)$$

Experiment Length

Parameter accuracy is dependent on the record length so that a sufficient amount of data points is available to give long term dynamics. The RLS parameters theoretically converge in $\cong (n+m)$ steps for a white noise input [15]. For the random step input sequence, 4 to 15 sample intervals were typically observed for $\hat{\theta}_i$ parameters to converge.

$$\{k\} \text{ length} \cong (4 \text{ to } 15) * h \quad (29)$$

Test Case

To verify the accuracy of the collected ID parameters, an analog ACSL [16] simulation of the R-C parameters for a single step input power pulse is done to compare the final estimated and actual TTIC.

A test case utilizing the calculated transistor thermal parameters of A1 & A2 was used to verify the TDT procedure. As a first step, the linear difference equation Eq(19) was used to calculate the true $y(k)$ rather than the actual TTIC. The final results are found in Section VI.

V. Parameter Estimation

Method

The goal of the identifier block in Fig. 3 is to determine unknown a_i and b_i filter coefficients of the parameter vector $\hat{\theta}$ in $\hat{\theta}_j c(q, \hat{\theta})$. Deriving a control law for parameter estimation from the parameter vector error $[\theta_i - \hat{\theta}_i]$ is not possible since $\hat{\theta}_i$ parameters are not directly measurable. However, a prediction estimate $\hat{y}(k)$ for every measurable / calculable $y(k)$ can be formulated. If a linear model is assumed, then the resulting prediction error estimate $[y(k) - \hat{y}(k)]$ is a function of θ as shown in Eq(19).

$$\text{error}(k, \hat{\theta}) = y(k) - \hat{y}(k) \quad (30)$$

Various methods that minimize the sum of the squares of this prediction error are Maximum Likelihood, Least Mean Squares, Extended Least Squares and Recursive Least Squares. The RLS method was chosen since it is computationally fast, requires no matrix inversions, and tends to converge rapidly. Observed convergence rates for the 2nd order model varied from 4 to 8 amplitude step changes, depending on the closeness of the $\hat{\theta}_i$ initial guesses. Thus, only a small number of $y(k)$ calculations using the TTIC are required. A disadvantage of RLS is the $\hat{\theta}_i$ "biasing" toward wrong values as $k \rightarrow \infty$ when the process output $y(k)$ is measured in the presence of noise. However, in this application, $y(k)$ is virtually noiseless since the primary source of noise appears to be interpolation errors when reading the TTIC curve.

Solution

A recursive identifier is one in which both input $u(k)$ from $\{k = 0$ to $N - 1\}$ and output $y(k)$ measurements from $\{k = 1$ to $N\}$ are made and inputed to the estimator to determine the parameter vector $\hat{\theta}$. The symbol N is the total number of sampling intervals over which the data is collected. The minimization process requires $N > n+m$ to effectively average out error residuals.

Defining the data regression vector as (31), then the error at any given time $k = \eta$ as a function of $\hat{\theta}$ is given by Eq(32).

$$\chi(k+1) = \begin{pmatrix} y(k) \\ \vdots \\ y(k+1-n) \\ u(k) \\ \vdots \\ u(k+1-m) \end{pmatrix} \quad (31) \quad \text{error}(\eta, \hat{\theta}) = y(\eta) - \chi^T(\eta) \hat{\theta} \quad (32)$$

The error vector equation and the error vector, output vector and regression vectors collected from time η to N is thus:

$$\varepsilon(N, \hat{\theta}) = Y(N) - \Psi(N) \hat{\theta} \quad (33)$$

where

$$\varepsilon(N, \hat{\theta}) = \begin{pmatrix} \varepsilon(\eta, \hat{\theta}) \\ \varepsilon(\eta+1, \hat{\theta}) \\ \vdots \\ \varepsilon(N, \hat{\theta}) \end{pmatrix} \quad Y(N) = \begin{pmatrix} y(\eta) \\ \vdots \\ y(\eta+1) \\ \vdots \\ y(N) \end{pmatrix} \quad \Psi(N) = \begin{pmatrix} \chi^T(\eta) \\ \chi^T(\eta+1) \\ \vdots \\ \chi^T(N) \end{pmatrix}$$

$$\varepsilon \in \mathbb{R}^{N-\eta+1} \quad Y \in \mathbb{R}^{(N-\eta+1) \times (n+m+1)} \quad \Psi \in \mathbb{R}^{(N-\eta+1) \times (n+m+1)}$$

The optimal $\hat{\theta}$ is the one that minimizes Eq(33) error in a least square sense. This requires a performance index $J(N, \hat{\theta})$ taking the gradient $\partial J / \partial \hat{\theta}$ and setting it equal to zero. This results in the well known least square solution of Eq(35). [15]

$$J(N, \hat{\theta}) = \sum_{k=\eta}^N \varepsilon(k, \hat{\theta}) \varepsilon(k, \hat{\theta}) \quad (34)$$

$$\hat{\theta}[N] = [\Psi(N) \Psi(N)]^{-1} \Psi^T(N) Y(N) = P[N] Y[N] \quad (35)$$

$$\hat{\theta}[N+1] = P[N+1] Y[N+1] \quad (36)$$

The optimal coefficients $\hat{\theta}$ will exist if the pseudoinverse $P[N]$ is nonsingular. This condition is satisfied assuming persistent excitation, which is guaranteed by selecting the amplitude of $u(k)$ randomly for times up to $k = N$ by using a pseudo random sequence. As new data arrives $\{u(N), y(N+1)\}$, the objective of RLS is to update $\hat{\theta}[N]$ to $\hat{\theta}[N+1]$ in terms of the old data and $\hat{\theta}[N]$ vector and similarly update $P[N]$ to $P[N+1]$. The $P[N+1]$ matrix in Eq(36) can be related to $P[N]$ without inversion via the matrix inversion lemma [17]. Thus $\hat{\theta}[N+1]$ can be related to $\hat{\theta}[N]$ without inversion by

$$\hat{\theta}(k+1) = \hat{\theta}(k) + L(k+1) [y(k+1) - \chi^T(k+1) \hat{\theta}(k)] \quad (37)$$

new estimate = old estimate + correction term

where $L(k+1)$ = Gain Matrix
 $y(k+1)$ = New Data Measured
 $\chi^T(k+1) \hat{\theta}(k)$ = Prediction of $y(k+1)$
 correction term = Gain * (y Output Error Equation)

Starting Conditions

A first estimate for $\hat{\theta}[N]$ without inversion may be obtained using the a-priori knowledge of Section III Parameter Initialization Procedure. Alternatively, Soderstrom's suggested starting conditions for $\hat{\theta}[N]$ and $P[N]$ can be utilized. Two initial conditions for $y(k)$ must be calculated for the 2nd order model, thus starting the process at $k = 2$.

$$\hat{\theta}[N] = 0 \quad P[N] = \alpha * I \quad \alpha = (10/n) \sum_{k=0}^1 y^2(k) \quad (38)$$

Algorithm

The identification scheme [18] used for the computer program is:

- (0) Initialize $\hat{\theta}[N]$ and $P[N]$; set $k = 2$
- (1) Form $\chi(k+1)$ data vector $= (n+m+1) * 1$ matrix for SISO system
- (2) Update $L(k+1) = [1/\gamma] P(k) \chi(k+1) [1/\alpha + \{ \chi^T(k+1) / \gamma \} P(k) \chi(k+1)]^{-1}$ where $\gamma = 1$
 $\alpha = 1$
 $\{ \ }^{-1}$ = simple inverse of a scalar value
 $L(k+1)$ = $(n+m+1) \times 1$ matrix
 $P(k)$ = $(n+m+1) \times (n+m+1)$ matrix
 $\chi(k+1)$ = $(n+m+1) \times 1$ vector
 $\chi^T(k+1)$ = $1 \times (n+m+1)$ vector
- (3) Measure $y(k+1), u(k+1)$
- (4) Update $\hat{\theta}(k+1) = \hat{\theta}(k) + L(k+1) [y(k+1) - \chi^T(k+1) \hat{\theta}(k)]$
- (5) Update $P(k+1) = [1/\gamma] \{ I - L(k+1) \chi^T(k+1) \} P(k)$
- (6) Replace k by $k+1$ and Go To (1)

VI. Results

$y(k)$ Calculation

Calculation of the true junction-case temperature, $\Delta T_{jc} = y(k)$, in the RLS routine must be done using the TTIC and a persistently excited device input power sequence, $u(k)$. The "Calculated Parameter" TTIC shown in Fig. 6 was specifically chosen as a test case example since the $R-C$ parameters generating this curve are exactly known and can be compared to "Identified" $R-C$ parameters. A computer program was written to calculate $y(k)$ utilizing input data points (10 pts / decade) from the test case TTIC curve. The program calculates ΔT_{jc} for a continuous $u(k)$ input sequence using the superposition of equivalent heating and cooling pulses.

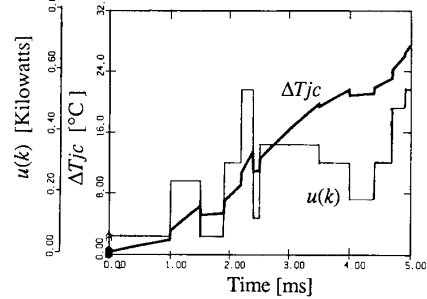


Fig. 8 Simulated Analog ΔT_{jc} Response to $u(k)$ Power Sequence

Fig. 8 shows a typical $u(k)$ input power sequence and output ΔT_{jc} response for a 4th order ACSL analog simulation model using the $R-C$ Calculated Parameters in Table 1. The basic $u(k)$ pulse pattern from $t = 0$ to $t = 2.5$ ms was repeated starting at $t = 2.5$ ms. Of particular interest is the instantaneous temperature jump at each new $u(k)$ pulse due to the P^*R1 component of Fig. 1c.

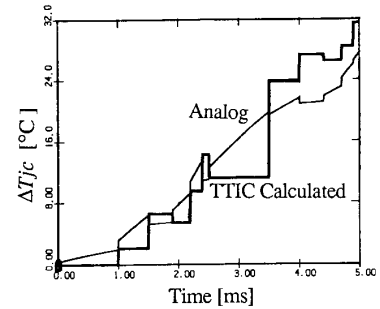


Fig. 9 TTIC Calculated & Analog ΔT_{jc} Responses to $u(k)$

Fig. 9 shows the program TTIC Calculated ΔT_{jc} discrete step response as well as the analog ΔT_{jc} response of Fig. 8. The TTIC inherently incorporates the heating / cooling "integration step" while providing "sampled" results at the end of each pulse. The discrete TTIC and analog Simulated ΔT_{jc} responses are virtually identical at the end of each pulse from $t = 0$ to $t = 2.5$ ms. At time $t = 2.5$ ms, the

Table 2 Summary of RLS Identification Runs

ID Run #	h (ms)	τ_{calc} (ms)	Initial Parameters									Identified Parameters									
			R1	C1	R2	C2	R3	C3	R4	C4	R5	R1	C1	R2	C2	R3	C3	R4	C4	R5	
Fast	1a	0.01	0.32	0.0000	0.0000	0.0000	0.0000	0.0000	-----	-----	-----	-----	0.0064	0.0324	0.0179	0.0005	0.0762	-----	-----	-----	
	1b	0.01	0.32	0.0080	0.3600	0.0220	0.4880	0.0780	-----	-----	-----	-----	0.0064	0.0329	0.1095	0.0727	-0.015	-----	-----	-----	
	2a	1.0	16	0.0064	0.0329	0.1095	0.0727	-0.015	-----	-----	-----	-----	0.0064	0.0329	0.1083	0.1697	0.1095	-----	-----	-----	
	2b	1.0	16	0.0064	0.0329	0.1095	0.0727	-0.015	-----	-----	-----	-----	0.0064	0.0330	0.1109	0.1479	0.1219	-----	-----	-----	
Slow	3	10	80	-----	-----	-----	-----	0.2397	0.6650	0.2910	0.9000	0.011	-----	-----	-----	0.2397	1.2263	0.0003	-0.117	0.1804	
	4	100	400	-----	-----	-----	-----	0.2397	1.2263	0.2910	0.9000	0.011	-----	-----	-----	0.2396	1.1799	0.1659	9.4842	0.0110	
True Value			-----	-----	-----	-----	-----	-----	-----	-----	-----	-----	0.0064	0.0330	0.111	0.1480	0.1220	1.1800	0.1660	9.5000	0.0110

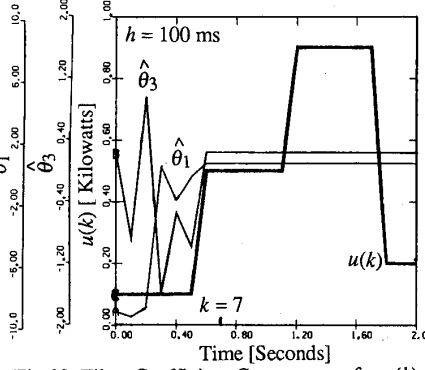
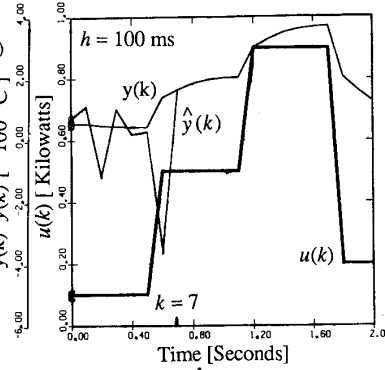
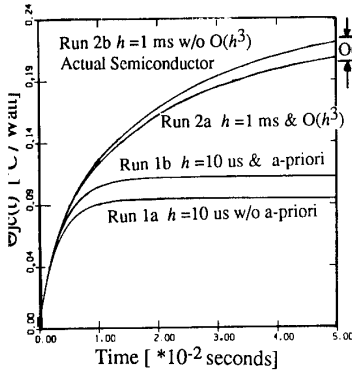


Fig. 10 Reconstructed vs. Actual TTIC's Fig. 11 $y(k)$ and $\hat{y}(k)$ Response to $u(k)$ Fig.12 Filter Coefficient Convergence for $u(k)$

TTIC program was restarted, retaining the ΔT_{jc} (2.5ms) value as a new starting point, and assuming the next pulse from $t = 2.5$ to 3.5 ms is of 1 ms duration. The resulting error that follows between the two ΔT_{jc} responses illustrates a typical misapplication of the TTIC concept that violates the single pulse - zero initial condition assumptions upon which the curve is based. This is further clarified by calculating both analog and TTIC program ΔT_{jc} responses at $t = 2.5$ ms and $t = 3.5$ ms.

The TTIC ΔT_{jc} response is calculated for a single equivalent heating pulse of 2.5 ms duration and an amplitude (165 watts) corresponding to the average value of $u(k)$ from $t = 0$ to $t = 2.5$ ms. The resulting temperature using the "Calculated Parameter" curve of Fig. 6 is 10.6°C and is in agreement with the actual 11.1°C value obtained with both the TTIC program and analog simulation.

The calculated TTIC ΔT_{jc} response at 3.5 ms cannot be done by restarting at $t = 2.5$ ms as described above. The correct method must assume the device power profile of Fig. 4a with $P1=165$ W, $P2 = 360$ W, $t1 = 2.5$ ms, and $t2 = 3.5$ ms. Using the TTIC of Figure 6 and Eq(3) results in less than 0.5°C error from the true temperature.

True RLS $y(k)$ calculation for the test case example with known $R-C$ parameters was done using Eq(19) discrete difference equation since it was easier to use and provided discrete temperature information identical to the TTIC computer program.

Test Case Results

Table 2 results show that only 4 basic identification runs were needed to identify the 9 unknown $R-C$ parameters. Correctly Identified Parameters in Table 2 are enclosed in a solid box. The four basic time constants listed in the τ_{calc} column are a direct result of the $\hat{\theta}$ Parameter Initialization Procedure using Eq(8e) thru (8i) and the Actual TTIC of Fig. 6. A suitable sampling rate for each ID run was derived from the τ_{calc} column by applying Eq(28).

The $R1, R2, R3, C1$ and $C2$ values were identified in 2 ID runs using the Fast 2nd Order $\hat{\theta}_{jc}$ model. Run 1a used the Soderstrom starting

conditions assuming no a-priori parameter information. As seen from Table 2, $R1$ and $C1$ are properly identified, as expected, for $h = 10$ usec. Run 1b shows that additionally $R2$ can be correctly identified by using all the a-priori Initial Parameter Estimates from Table 1. Following the TDT procedure, the output Identified Parameters from Run 1b were used as Initial Parameter estimates for Run 2 using $h = 1.0$ ms. Run 2a results show that the final $R3$ and $C2$ values enclosed by the dashed box in Table 2 are within 20% of the actual values. This error is due to the method of calculating true $y(k)$ using the linear difference equation rather than the TTIC approach. This is caused by the DC gain discretization error introduced by the $O(h^3)$ term in the b_2 filter coefficient of Appendix 2. Run 2b removed this error term resulting in final Identified Parameters within 0.1 % accuracy. Figure 10 shows a graphical comparison of ID Runs 1a, 1b, 2a, and 2b by reconstructing the TTIC from corresponding identified $R-C$ parameters.

Components $R3, R4, R5, C3$ and $C4$ were identified using the slow $\hat{\theta}_{jc}$ model in Runs 3 and 4. The Initial Parameter estimate for $R3$ in Run 3 was determined by adding the $R1, R2$ and $R3$ values identified from Run 2b. The remaining Initial Parameters were obtained from the $\hat{\theta}$ Initialization Procedure and Table 1. The results from Run 3 show that, similar to Run 1b, the $R3-C3$ parameters associated with the faster time constant are correctly identified. Run 4 used these two component values along with estimates for $R4, R5$ and $C4$ from Table 1 as input parameters. The final Identified Parameters were within 0.2% of the actual values. Figures 11 and 12 show typical identification waveforms for Run 4 with $h = 100$ ms.

Fig. 11 illustrates the output error equation, $\tilde{y}(k) = y(k) - \hat{y}(k)$, being driven to zero in $k = 7$ samples. After time $k = 7$, the $y(k)$ and $\hat{y}(k)$ temperature response waveforms to the $u(k)$ sequence are identical. The $u(k)$ power sequence consists of six similar amplitude pulses before changing at $k = 7$. Figure 12 depicts the same $u(k)$ power sequence applied as in Fig.11. In addition, a typical RLS convergence pattern for two $\hat{\theta}$ parameter values is shown. The digital filter coefficients, $\hat{\theta}_1 = a_1$ and $\hat{\theta}_3 = b_0$, are shown dynamically adapting to new values to satisfy the $\tilde{y}(k)$ error equation and also reach steady state at $k = 7$ samples. The steady state values of the five identified $\hat{\theta}_i$ were used to further extract the model $R-C$ parameters.

VIII. Conclusion

This paper has identified a need for a device thermal model to maximize simulation capability by solving both the electrical and thermal network models simultaneously. A new approach to the problem of determining device thermal model characteristics was presented using System Identification concepts. Governing thermal laws, device physical packaging construction, manufacturer's data specification sheet and standard TTIC "graphical transfer function" information were used as a-priori knowledge to determine the model order and structure. The typical semiconductor model structure inherently contains wide time scale separation of the thermal time constants. This information was used to advantage in formulating a new "systematic" thermal R-C Extraction procedure using RLS and time decoupled theory. Time decoupled theory used multiple RLS identification runs to estimate the major time constants in succession from the fastest to the slowest. The identified R-C parameters from each run are found from the digital filter coefficients of the estimated "mathematical" transfer function Θ_{jc} . A test example using a transistor module was used to verify the proposed technique. Calculated R-C parameters obtained from physical dimensions were also performed.

The proposed parameter identification concept may also be extended to other thermal systems with inherent overload capability such as transformers, rotating machines, etc. A TTIC curve can be generated for a step input of equipment power using output temperature data or possibly internal temperature states. However, a linearized model for a range of input power will be obtained due to nonlinear convection and radiation heat transfer. Time decoupled theory may possibly be extended to rotor time constant identification in AC vector control.

Appendix A1 Governing Thermal Laws

A R-C parameter model of Fig. 1c and Fig. 5d is assumed for 50 Amp, 500 Volt, 300 Watt dual darlington isolated base transistor module with base dimensions of 1.25 x 3.6 inches. The specified $\Theta_{jc}(tss) = 0.41$ [°C / W] and has the TTIC curve shown in Fig. 6. Thermal resistance is directly analogous to Ohm's law for electrical resistance. A 45° angle from the junction to the case was assumed

$$R_t = L_t / (K * Ae) \quad (A1)$$

L_t = thickness[inch] of material along heat flow path

K = material thermal conductivity from Table A2

Ae = cross section area [sq.in.] perpendicular to heat flow path in calculating an effective heat spreader area $\{Ae\}$ for succeeding layers with a much greater true cross sectional area $\{Ac\}$. Thermal resistance $\{R_i\}$ is calculated to the midpoint of each major heat capacity material where the capacitance $\{C_i\}$ is assumed a lumped parameter. However, the silicon chip is an exception where it is assumed that the top 1/2 of the thickness $\{L_t\}$ is really the distributed power source $\{P\}$. The ceramic insulator (26%) and solder interfaces (34%) account for 50% of the specified $\Theta_{jc}(tss)$.

The thermal capacitance is calculated using true material volume:

$$C_i = \rho * Cp * V \quad (A2)$$

ρ = density of the material from Table A2

Cp = specific heat of the material from Table A2

V = true material volume $(L_t * Ac)$ [in³]

Table A1 Calculated Parameter Spreadsheet

Material	t (mils)	Ac (sq. inch)	Ae (sq. inch)	Li	Rind (C/W)	Rt (C/W)	Ci (W-s/C)
Silicon	15	0.43 x 0.43	0.43 x 0.43	= t/4 = t/4	0.008 0.008	0.008	0.033
Sn-Pb 10-90	3	0.43 x 0.43	0.43 x 0.43	= t	0.087	0.111	
Mo	20	0.43 x 0.43	0.43 x 0.43	= t/2 = t/2	0.016 0.016	0.148	
Sn-Pb 75-25	17	0.43 x 0.43	0.43 x 0.43	= t	0.097	0.122	
Cu block	63	0.6 x 0.6	0.6 x 0.6	= t/2 = t/2	0.0089 0.0089	1.18	
Sn-Pb 75-25	4	0.6 x 0.6	0.6 x 0.6	= t	0.012		
Al2O3	20	0.6 x 0.6	0.6 x 0.6	= t	0.109	0.166	
Sn-Pb 75-25	9	0.6 x 0.6	0.6 x 0.6	= t	0.027		
Cu base	118	1.2 x 1.2	0.75x0.75	= t/2 = t/2	0.010 0.010	0.010	9.50

Table A2 Material Properties Assumed

Material	K[W/(°C-in)]	p[lb/cu.in]	Cp[W-S/lb-°C]
Silicon	2.134	0.083	303
280°C Solder (Sn10-Pb90)	0.914	-----	----
Molydendum	3.296	0.369	115
180°C Solder (Sn75-Pb25)	0.914	-----	----
Ceramic	0.510	-----	----
Copper	9.77	0.320	175

Appendix A2 Discrete Transfer Function Model

Let the following constants be defined for a chosen sample time h :

$$K_{11} = h/R_1C_1 \quad K_{21} = h/R_2C_1 \quad K_{22} = h/R_2C_2 \quad K_{32} = h/R_3C_2$$

$$\Phi = \begin{pmatrix} (1-K_{11}) & K_{21} \\ K_{22} & (1-K_{22}-K_{32}) \end{pmatrix} \quad \Gamma = (h/C_1) * \begin{pmatrix} (1-0.5 * K_{21}) \\ 0.5 * K_{22} \end{pmatrix}$$

$$C = [1 \quad 0] \quad D = [R_1]$$

The digital filter coefficients used in Eq(18) transfer function are:

$$a_1 = -2.0 + K_{21} + K_{22} + K_{32}$$

$$a_2 = 1.0 - K_{21} - K_{22} - K_{32} + K_{21} * K_{32} \quad ; \quad b_0 = R_1$$

$$b_1 = R_1 \{ -2.0 + K_{11} + K_{21} + K_{22} + K_{32} - 0.5 * K_{11} * K_{21} \}$$

$$b_2 = R_1 \{ 1.0 - K_{11} - K_{21} - K_{22} - K_{32} + 0.5 * K_{11} * K_{21} + \dots$$

$$K_{11} * K_{22} + K_{11} * K_{32} + K_{21} * K_{32} - O(h^3) \}$$

where the $O(h^3) = 2nd$ order error term = $0.5 * K_{11} * K_{21} * K_{32}$

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